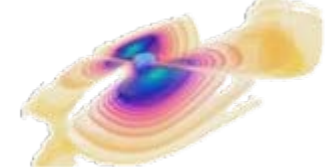


Iteration stability of the Einstein constraints in the conformal thin-sandwich approach

Charalampos Markakis (FSU Jena)
Richard H Price (UT Brownsville)
John L Friedman (UW Milwaukee)
Bernd Brügmann (FSU Jena)

NEB 15
Chania
June 23, 2012



Iteration stability - Applications

- Single, non rotating star

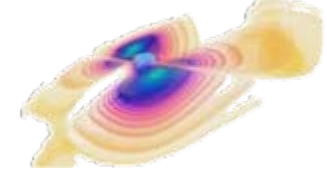


- Single, rotating star



- Binary system





Articles where methods are used

- R. James, *Astrophys. J.* 140, 552 (1964)
- R. Stoeckly, *Astrophys. J.* 142, 208 (1965)
- J. P. Ostriker and J. W.-K. Mark, *Astrophys. J.* 151, 1075 (1968)
- I. Hachisu, *Astrophys. J. Suppl.* 61, 479 (1986)
- E. Butterworth and J. Ipser, *Astrophys. J.* 204, 200 (1976)
- M. Ansorg, A. Kleinwächter, and R. Meinel, *Astr. Astrophys.* 381, L49 (2002)
- S. Bonazzola and S. Schneider, *Astrophys. J.* 191, 195 (1974)
- S. Bonazzola, E. Gourgoulhon, M. Salgado, and J. Marck, *Astron. Astrophys.* 278, 421 (1993)
- H. Komatsu, Y. Eriguchi, and I. Hachisu, *Mon. Not. R. Astron. Soc.* 237, 355 (1989)
- H. Komatsu, Y. Eriguchi, and I. Hachisu, *Mon. Not. R. Astron. Soc.* 239, 153 (1989)
- N. Stergioulas, *Living Reviews in Relativity* (2003)
- U. M. Schaudt, *Ann. Henri Poincaré* 1 (2000) 945 – 976
- U. M. Schaudt & H. Pfister, *Phys. Rev. Lett.* 77, 16 (1996)
- ...

Equations of Stellar Structure



Poisson

$$\Phi(\vec{r}) = -G \int d^3r' \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|}$$

Euler

$$\nabla(h + \Phi - \frac{1}{2}v^2) = 0$$

EOS

$$h := \int_0^\rho \frac{dP}{\rho}$$

Equations of Stellar Structure



Poisson

$$\Phi(\vec{r}) = -G \int d^3r' \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|}$$

Euler

$$h + \Phi - \frac{1}{2} v^2 = \kappa$$

EOS

$$\rho = \rho[h]$$

Iterative procedure (SCF method)



Poisson

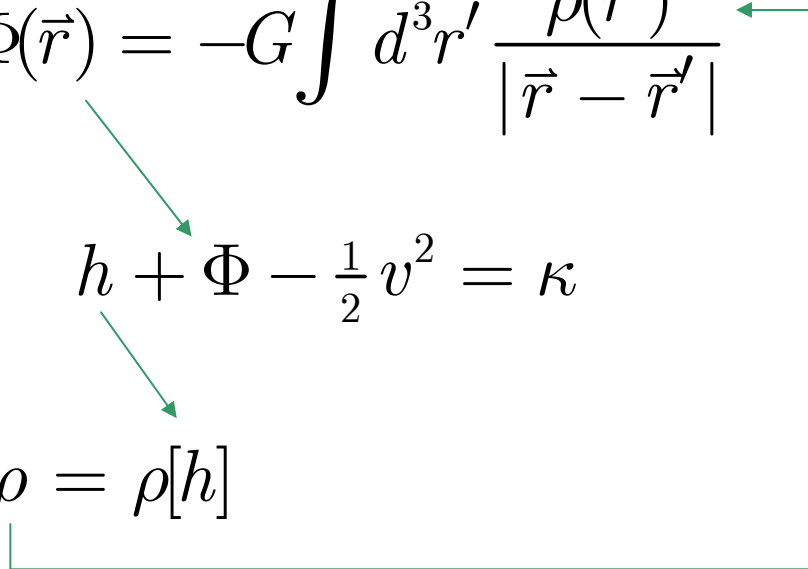
$$\Phi(\vec{r}) = -G \int d^3r' \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|}$$

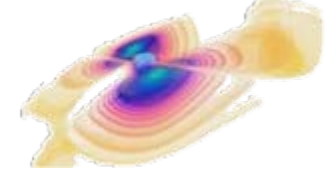
Euler

$$h + \Phi - \frac{1}{2} v^2 = \kappa$$

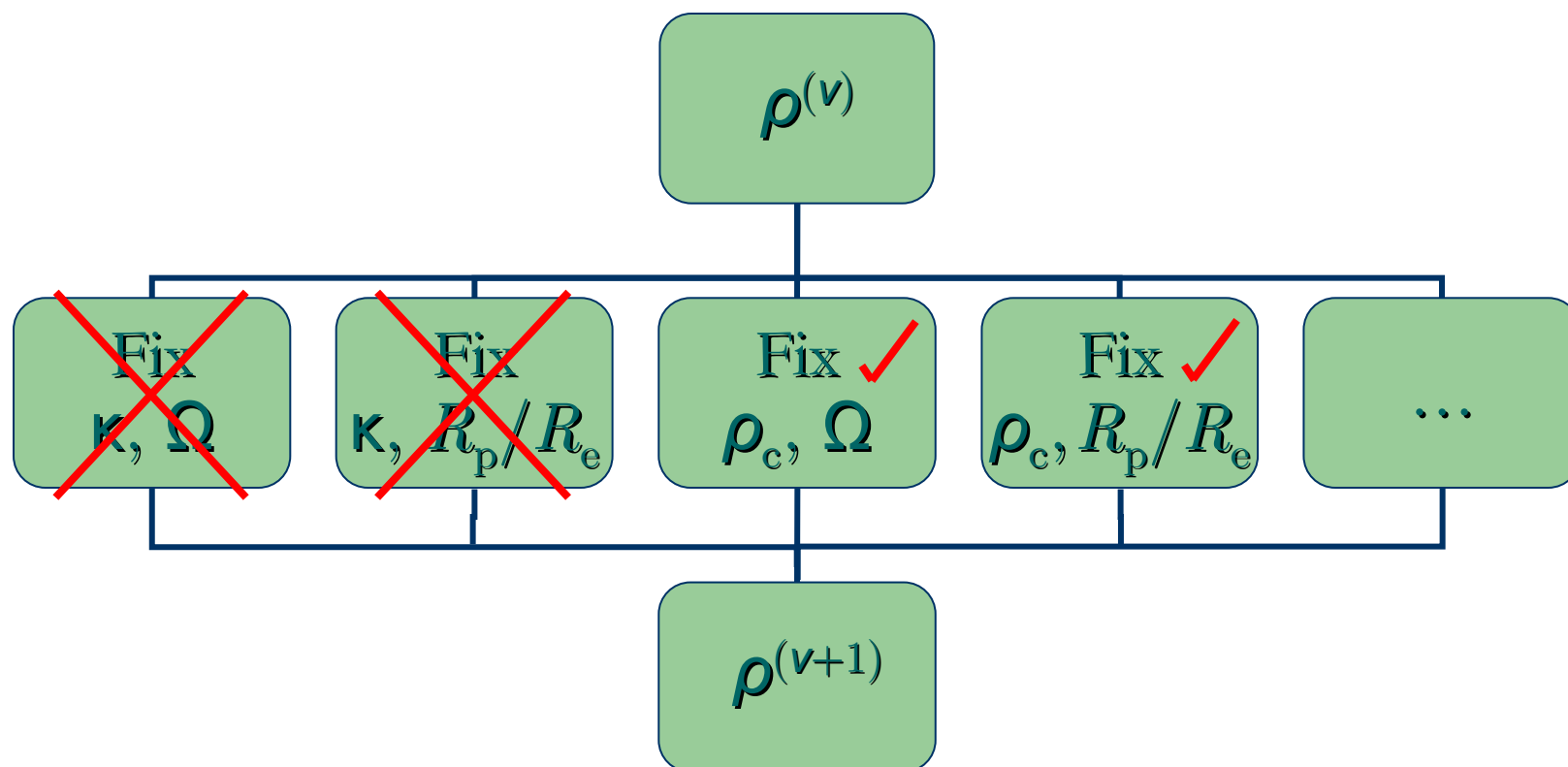
EOS

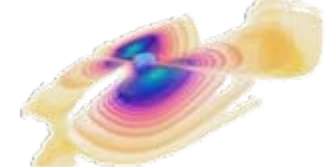
$$\rho = \rho[h]$$



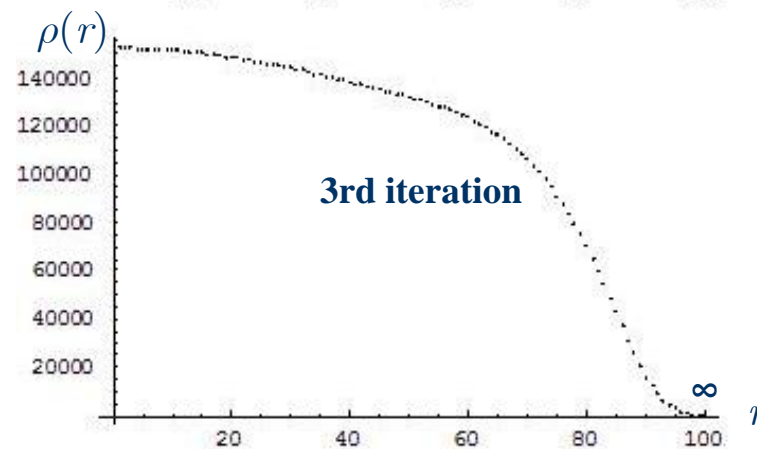
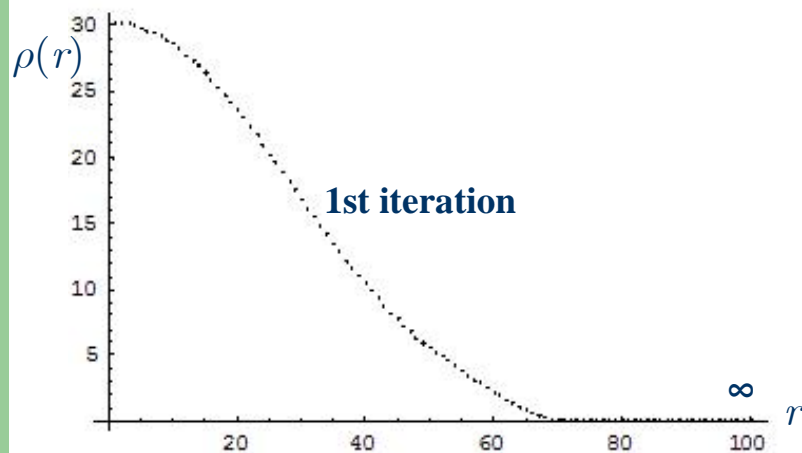
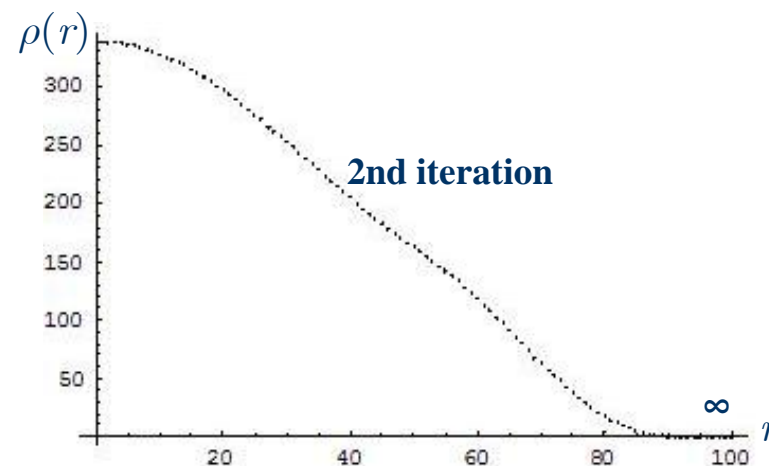
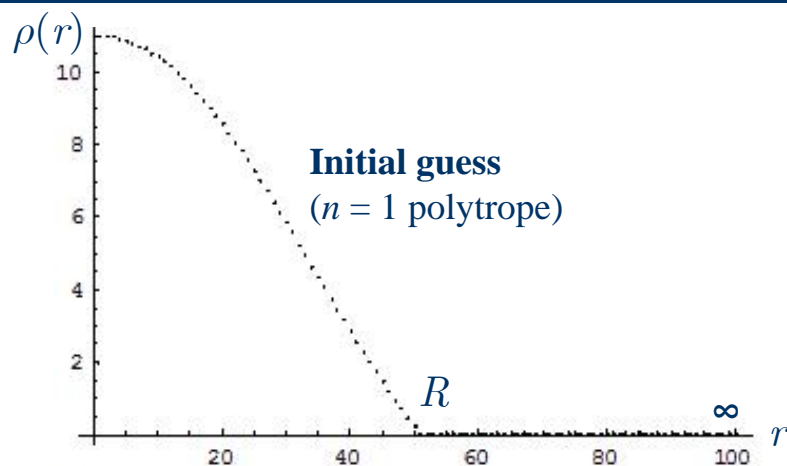


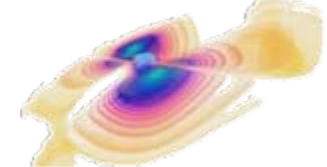
Variants of the SCF method



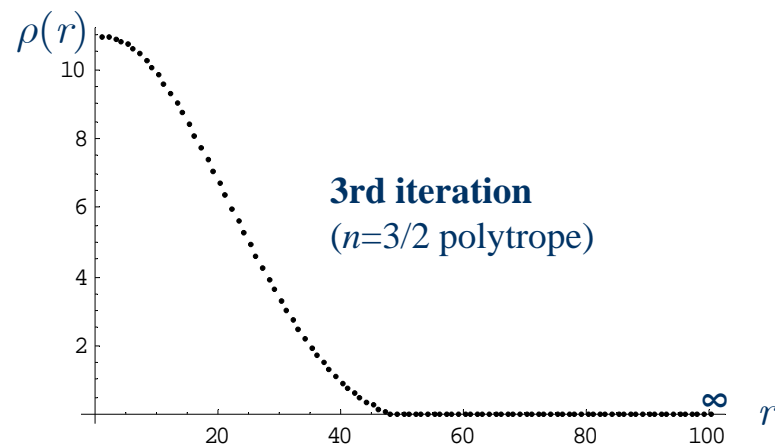
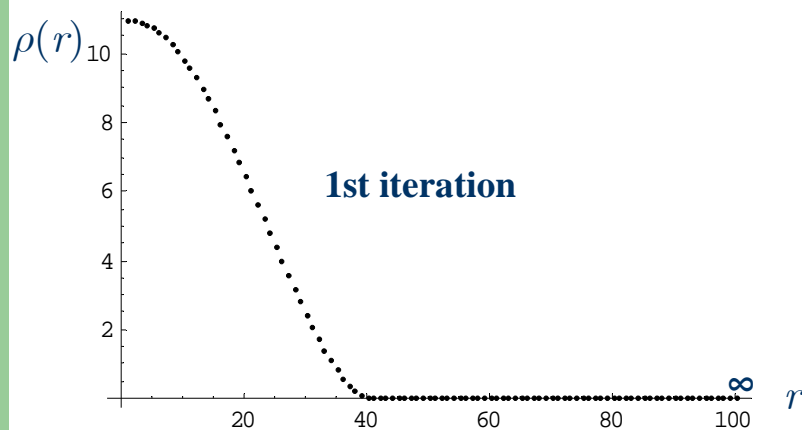
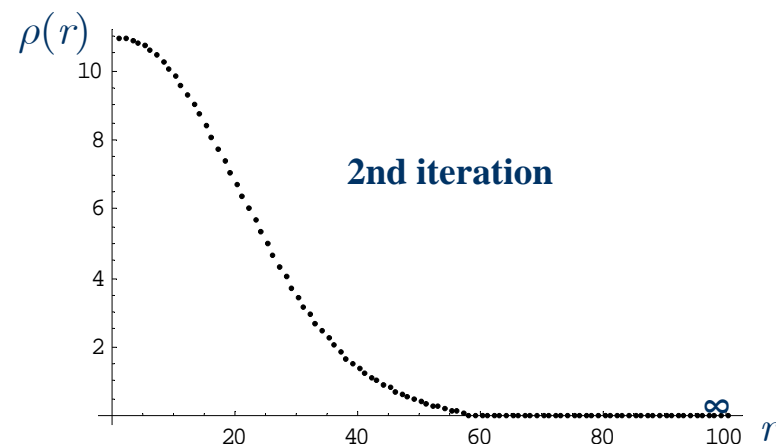
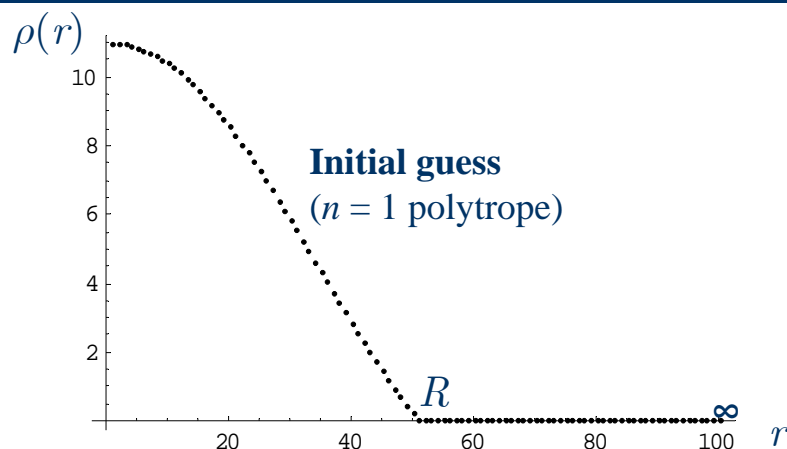


Iteration with fixed κ





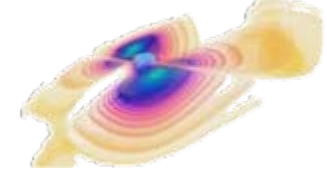
Iteration with fixed central density



Iteration Stability



- SCF method has been used for > 40 years
- Little is known about why it works when one quantity is fixed and not another...
- In the past, obtaining numerical solutions with SCF was too easy for people to care why it worked...
- Not the case for relativistic binaries, so understanding convergence is crucial today...
- Consider behavior of iterative scheme near an exact solution



Iteration Stability

Exact solution

$$\Phi(\vec{r}) = -\frac{1}{4\pi} \int d^3r' \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|}$$

$$\eta + \Phi - \frac{1}{2}v^2 = \kappa$$

$$\rho = \eta^n$$

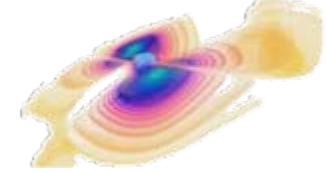
Linear perturbation

$$\delta\Phi(\vec{r}) = -\frac{1}{4\pi} \int d^3r' \frac{\delta\rho(\vec{r}')}{|\vec{r} - \vec{r}'|}$$

$$\delta\eta + \delta\Phi = \delta\kappa$$

$$\delta\rho = n\eta^{n-1}\delta\eta$$

Iteration operator $\eta^{\text{new}}(\vec{r}) = \kappa + \frac{1}{4\pi} \int d^3r' \frac{[\eta^{\text{old}}(\vec{r}')]^n}{|\vec{r} - \vec{r}'|}$



Iteration Stability

Exact solution

$$\Phi(\vec{r}) = -\frac{1}{4\pi} \int d^3r' \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|}$$

$$\eta + \Phi - \frac{1}{2}v^2 = \kappa$$

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Linear perturbation

$$\delta\Phi(\vec{r}) = -\frac{1}{4\pi} \int d^3r' \frac{\delta\rho(\vec{r}')}{|\vec{r} - \vec{r}'|}$$

$$\delta\eta + \delta\Phi = \delta\kappa$$

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Linear iteration operator $\delta\eta^{\text{new}}(\vec{r}) = \delta\kappa + \frac{n}{4\pi} \int d^3r' \eta^{n-1}(\vec{r}') \frac{\delta\eta^{\text{old}}(\vec{r}')}{|\vec{r} - \vec{r}'|}$



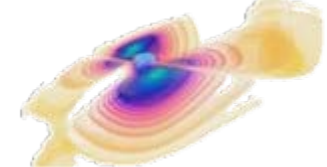
Iteration Stability

Linear iteration operator

$$\delta\eta^{\text{new}}(\vec{r}) = \delta\kappa + \frac{n}{4\pi} \int d^3r' \eta^{n-1}(\vec{r}') \frac{\delta\eta^{\text{old}}(\vec{r}')}{|\vec{r} - \vec{r}'|} \equiv L(\delta\eta^{\text{old}})$$

Fixed- κ iteration: $\delta\kappa = 0$

Fixed- $\eta(0)$ iteration: $\delta\eta(0) = 0$



Iteration Stability

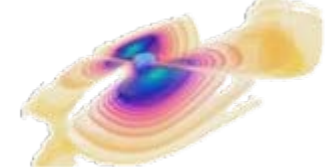
Linear iteration operator

$$\delta\eta^{\text{new}}(\vec{r}) = \delta\kappa + \frac{n}{4\pi} \int d^3r' [\eta(\vec{r}')]^{n-1} \frac{\delta\eta^{\text{old}}(\vec{r}')}{|\vec{r} - \vec{r}'|} \equiv L(\delta\eta^{\text{old}})$$

Eigenvalue problem

$$L(\delta\eta) = \lambda\delta\eta$$

Necessary condition for convergence: $|\lambda| < 1$



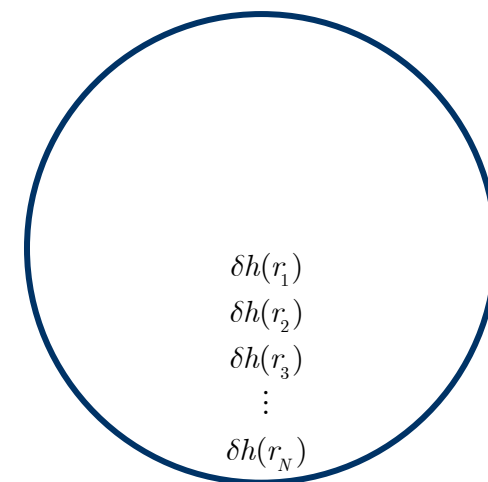
Iteration with fixed κ

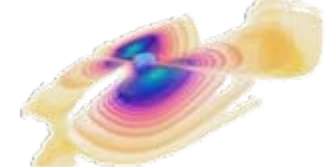
Linear iteration operator

$$\delta h^{\text{new}} = G \int d^3 r' \frac{d\rho}{dh} \frac{\delta h^{\text{old}}}{|\vec{r} - \vec{r}'|} \equiv L (\delta h^{\text{old}})$$

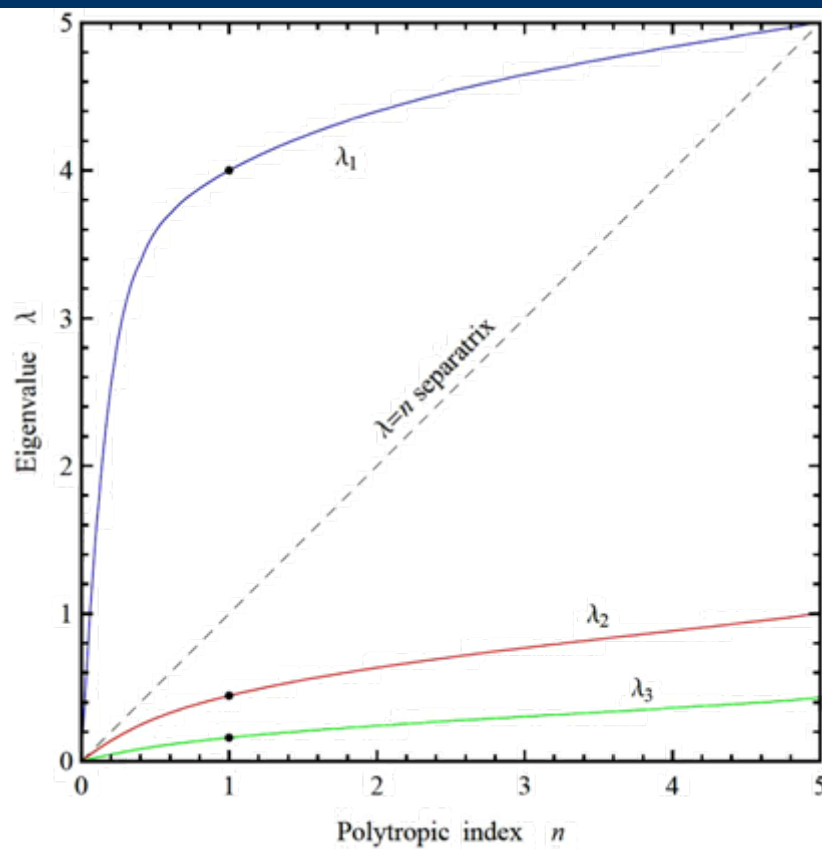
Discretization

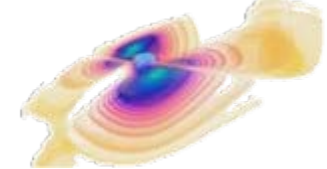
$$\delta h_i^{\text{new}} = \sum_j L_{ij} \delta h_j^{\text{old}}$$





Iteration with fixed κ



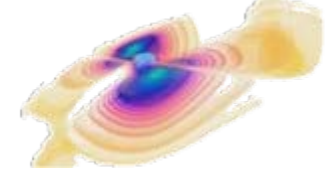


Iteration with fixed central density

$$\delta h^{\text{new}} = G \int d^3 r' \frac{d\rho}{dh} \left(\frac{1}{|\vec{r} - \vec{r}'|} - \frac{1}{|\vec{r}'|} \right) \delta h^{\text{old}} \equiv L(\delta h^{\text{old}})$$

Discretization

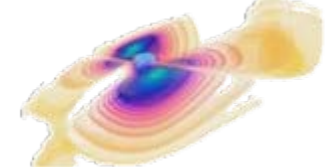
$$\delta h_i^{\text{new}} = \sum_j L_{ij} \delta h_j^{\text{old}}$$



Iteration with fixed central density

Jordan form of L

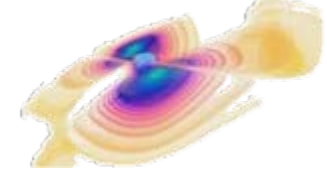
$$\mathbf{J} = \begin{pmatrix} 0 & 1 & & & \\ & 0 & 1 & & \\ & & 0 & \ddots & \\ & & & \ddots & 1 \\ & & & & 0 \end{pmatrix}$$



Iteration with fixed central density

Jordan form of L

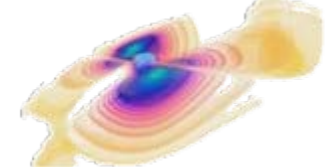
$$\mathbf{J}^2 = \begin{pmatrix} 0 & 0 & 1 & & \\ & 0 & 0 & \ddots & \\ & & 0 & \ddots & 1 \\ & & & \ddots & 0 \\ & & & & 0 \end{pmatrix}$$



Iteration with fixed central density

Jordan form of L

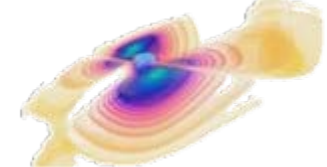
$$\mathbf{J}^{N-1} = \begin{pmatrix} 0 & \dots & 0 & 1 \\ & 0 & & 0 \\ & & 0 & \vdots \\ & & & \ddots \\ & & & & 0 \end{pmatrix}$$



Iteration with fixed central density

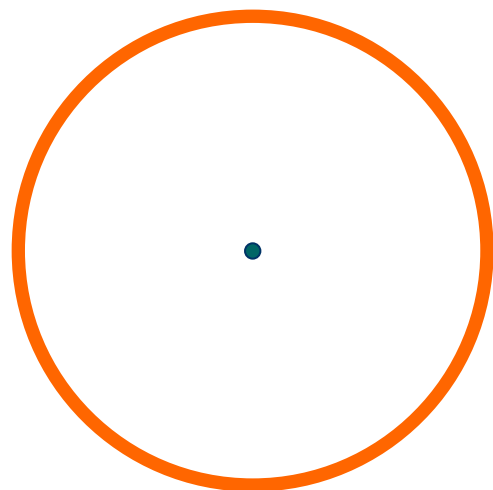
Jordan form of L

$$\mathbf{J}^N = \begin{pmatrix} 0 & \dots & 0 & 0 \\ & 0 & & 0 \\ & & 0 & \vdots \\ & & & \ddots \\ & & & & 0 \end{pmatrix}$$



Explanation

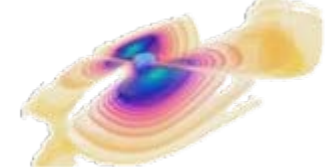
Fixed- ρ_0 linear iteration - explanation



$$\delta\Phi(\vec{r}) = -\frac{1}{4\pi} \int d^3r' \frac{\delta\rho(\vec{r}')}{|\vec{r} - \vec{r}'|}$$

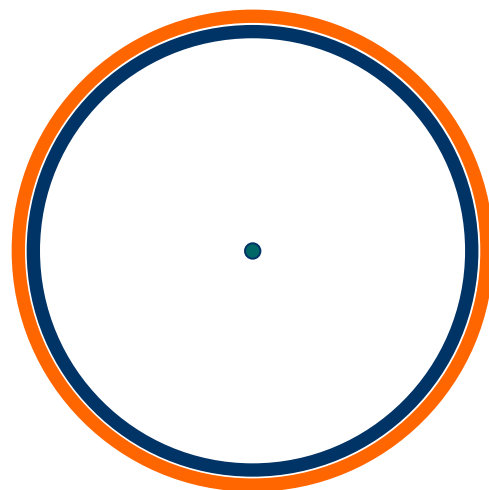
$$\delta\eta + \delta\Phi = \delta\kappa$$

$$\delta\rho = \frac{d\rho}{d\eta} \delta\eta$$



Explanation

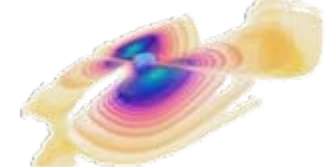
Fixed- ρ_0 linear iteration - explanation



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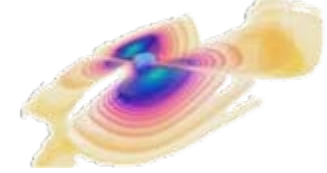
$$\delta\rho = \frac{d\rho}{d\eta} \delta\eta$$



Einstein constraints

Accurate calculation of physically relevant initial data corresponding to a system of black holes and/or neutron stars is an essential part of binary inspiral simulations.

In the 3+1 formalism, Einstein's equations are solved as a Cauchy-problem with initial data prescribed on a spacelike hypersurface. These initial data must satisfy the constraint equations. However, iterative algorithms for solving these equations may diverge, or may converge to an unphysical solution...



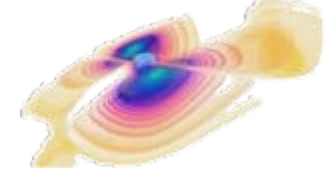
Euler – XCTS system

Hamiltonian constraint $\nabla^2 \Psi = -\frac{1}{2} \epsilon \Psi^5$

Maximal slicing condition $\nabla^2 (\alpha \Psi) = \frac{1}{2} (\epsilon + 6p) \alpha \Psi^5$

1st integral to Euler equation $\alpha h = \kappa$

EOS $\epsilon = \epsilon[h], \quad p = p[h]$



Euler – XCTS system

Hamiltonian constraint

$$\Psi(\vec{r}) = 1 - \frac{1}{8\pi} \int d^3r' \frac{1}{|\vec{r} - \vec{r}'|} \epsilon(\vec{r}') \Psi^5(\vec{r}')$$

Maximal slicing condition

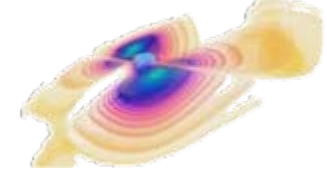
$$\alpha(\vec{r}) \Psi(\vec{r}) = 1 + \frac{1}{8\pi} \int d^3r' \frac{1}{|\vec{r} - \vec{r}'|} [\epsilon(\vec{r}') + 6p(\vec{r}')] \alpha(\vec{r}') \Psi^5(\vec{r}')$$

1st integral to Euler equation

$$\alpha h = \kappa$$

EOS

$$\epsilon = \epsilon[h], \quad p = p[h]$$



Euler – XCTS system

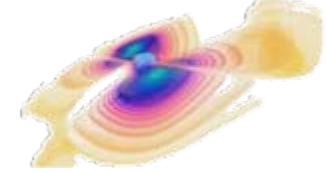
Hamiltonian constraint

$$\Psi(\vec{r}) = 1 - \frac{1}{8\pi} \int d^3r' \frac{1}{|\vec{r} - \vec{r}'|} \epsilon(\vec{r}') \Psi^5(\vec{r}')$$

Similar to updating of Newtonian enthalpy for fixed κ

$$\eta(\vec{r}) = \kappa + \frac{1}{4\pi} \int d^3r' \frac{1}{|\vec{r} - \vec{r}'|} \eta^n(\vec{r}')$$

Newtonian experience suggests eigenvalues > 1 ,
but $\epsilon(\vec{r})$ has compact support



Euler – XCTS system

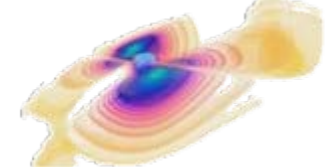
Hamiltonian constraint

$$\delta\Psi(\vec{r}) = -\frac{5}{8\pi} \int d^3r' \frac{1}{|\vec{r} - \vec{r}'|} \epsilon(\vec{r}') \Psi^4(\vec{r}') \delta\Psi(\vec{r}')$$

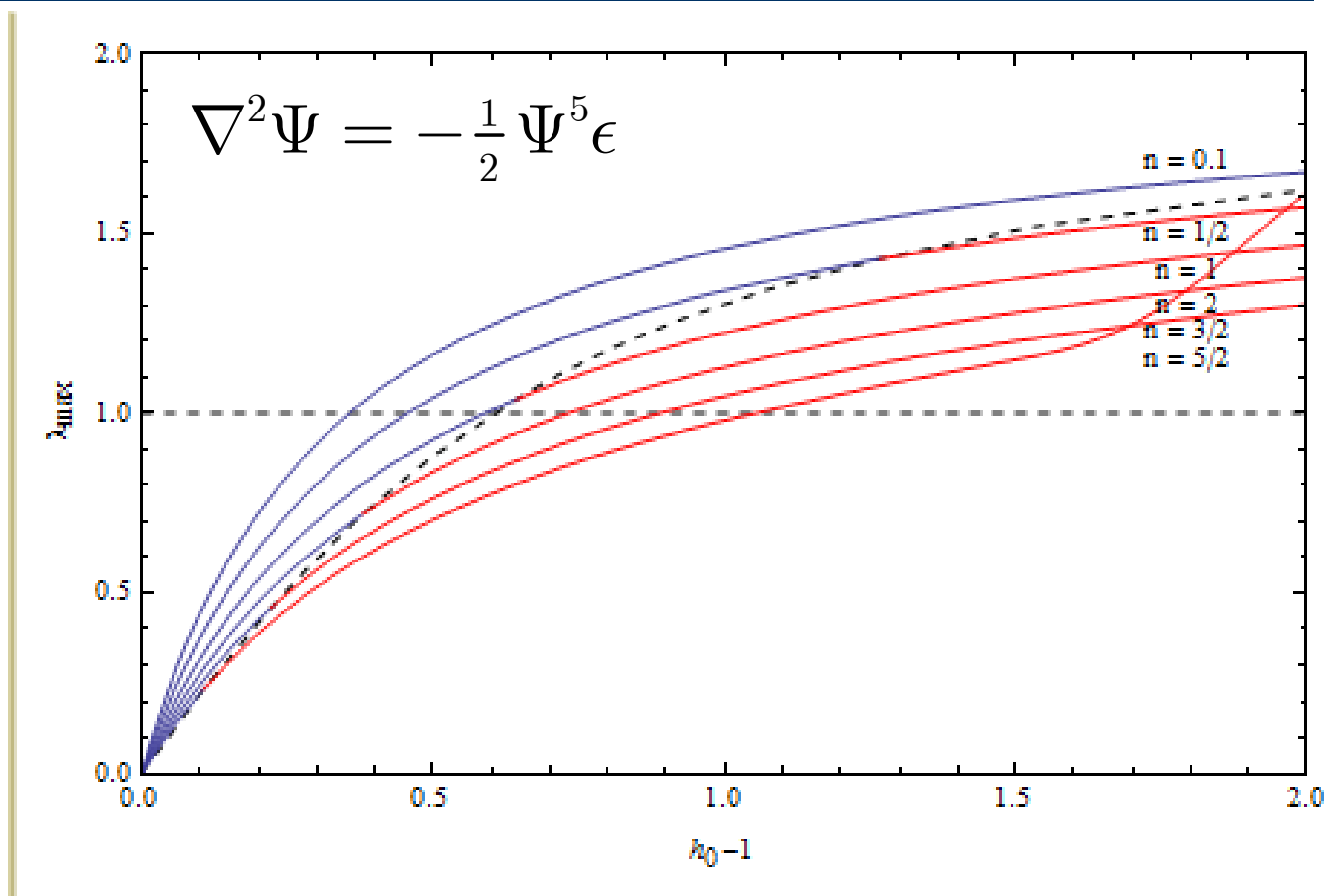
Similar to updating of Newtonian enthalpy for fixed κ

$$\delta\eta(\vec{r}) = \frac{n}{4\pi} \int d^3r' \frac{1}{|\vec{r} - \vec{r}'|} \eta^{n-1}(\vec{r}') \delta\eta(\vec{r}')$$

Newtonian experience suggests eigenvalues > 1 ,
but $\epsilon(\vec{r})$ has compact support



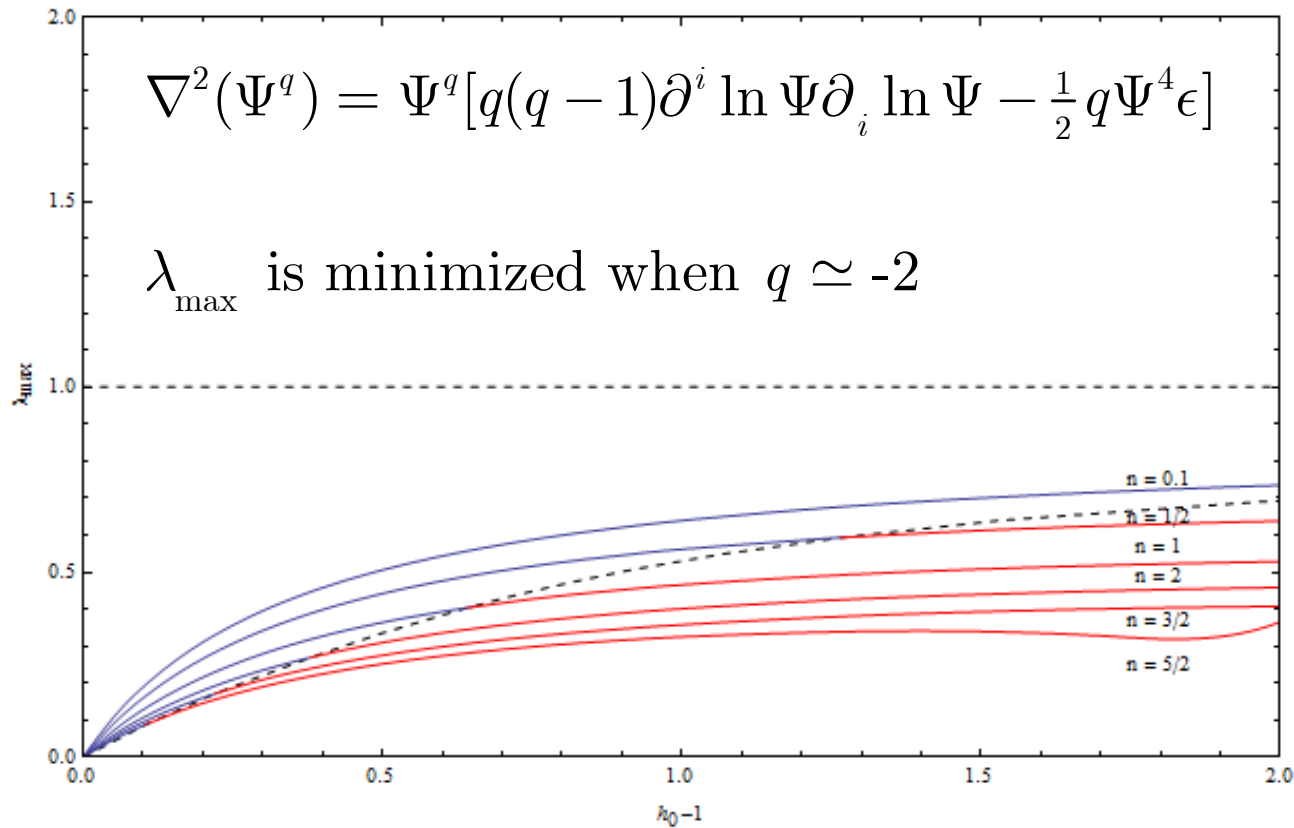
Hamiltonian constraint: iteration stability

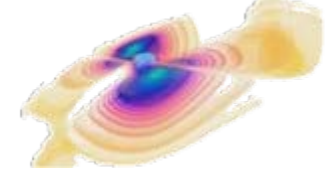


Hamiltonian constraint: iteration stability

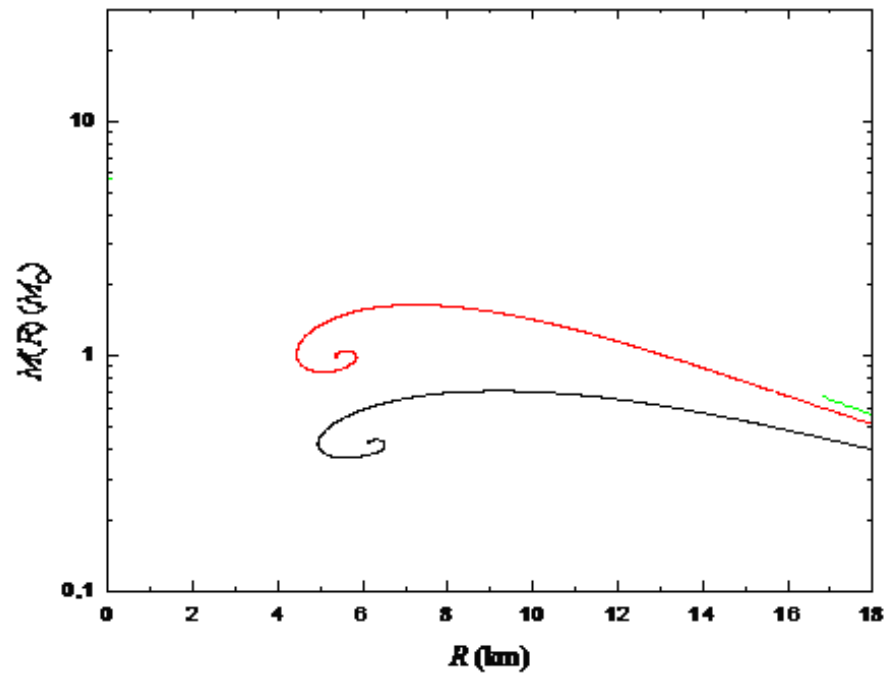
$$\nabla^2(\Psi^q) = \Psi^q \left[q(q-1) \partial^i \ln \Psi \partial_i \ln \Psi - \frac{1}{2} q \Psi^4 \epsilon \right]$$

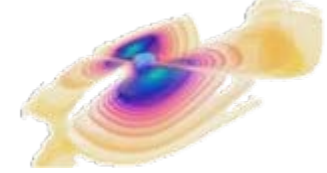
λ_{\max} is minimized when $q \simeq -2$





Turning point stability





Iteration Stability

Newtonian problem (completed):

R.H. Price, C. Markakis, J.L. Friedman

Iteration stability for simple Newtonian stellar systems

J. Math. Phys. 50, 073505 (2009)

Relativistic problem (in progress):

C. Markakis et al

Einstein constraints & fluid equilibria: optimizing convergence & stability of the extended conformal thin-sandwich system

- Study effect of rescaling source terms on stability
- Analyze full Euler-XCTS system
- Convergence optimization

